

# Temperature dependence of the critical field of the organic superconductor $\kappa$ -(BEDT-TTF)<sub>2</sub>Cu[N(CN)<sub>2</sub>]Br

V. A. Gasparov<sup>+1)</sup>, A. Audouard\*, L. Drigo<sup>×</sup>, J. A. Schlueter<sup>°</sup>

<sup>+</sup>*Institute of Solid State Physics Russian Academy of Sciences, 142432 Chernogolovka, Russia*

<sup>\*</sup>*Laboratoire National des Champs Magnetiques Intenses (UPR 3228 CNRS, INSA, UPS, UGA), F-31400 Toulouse, France*

<sup>×</sup>*Géosciences Environnement, 31400 Toulouse, France*

<sup>°</sup>*Materials Science Division, Argonne National Laboratory, Argonne, IL 60439-4845, USA*

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We present upper critical magnetic fields data for  $\kappa$ -(BEDTTTF)<sub>2</sub>Cu[N(CN)<sub>2</sub>]Br (hereafter k-Br), single crystals with magnetic field applied either perpendicular or parallel to the conducting layers. H-T phase diagram in this compound has been deduced from tunnel diode oscillator-based contactless measurements in pulsed magnetic fields up to 56 T for the inter-plane (H||b) and in-plane (H||ac) field directions. Temperature dependence of the magnetic penetration depth in DC fields is also reported. The temperature dependence of the upper critical magnetic field  $H_{c2}$  is not accounted for the Werthamer–Helfand–Hohenberg (WHH) model for in-plane configuration. For the inter-plane orientation, a significant upward curvature is observed as temperature decrease, in the range close to  $T_c$ , followed by saturation at lower temperatures. Possibly, this upturn is an indication of flux-line lattice melting.

The temperature-dependence of the upper critical magnetic field  $H_{c2}(T)$  in organic superconductors is most likely due to two independent pair-breaking mechanism [1, 2]: (i) close to  $T_c$ . Cooper pairing is suppressed by orbital currents that screen the external field, according to the well-known WHH model; (ii) towards lower temperatures, the limiting effect is caused by the Zeeman splitting, i.e., when the Zeeman energy becomes larger than the condensation energy, the Pauli limit,  $H_p$ , is reached. This paramagnetic limit,  $H_{c2}(T)$ , is lower than the orbital one which is related to the slope  $dH_{c2}(T)/dT$  close to  $T_c$ . Namely, when including Pauli paramagnetism, the upper critical field is reduced relatively to  $H_{c2}(T)$ :

$$H_{c2}(T) = H_{c2}^*(T) / \sqrt{1 + \alpha^2(T)}, \quad (1)$$

<sup>1)</sup>e-mail: vgasparov@issp.ac.ru

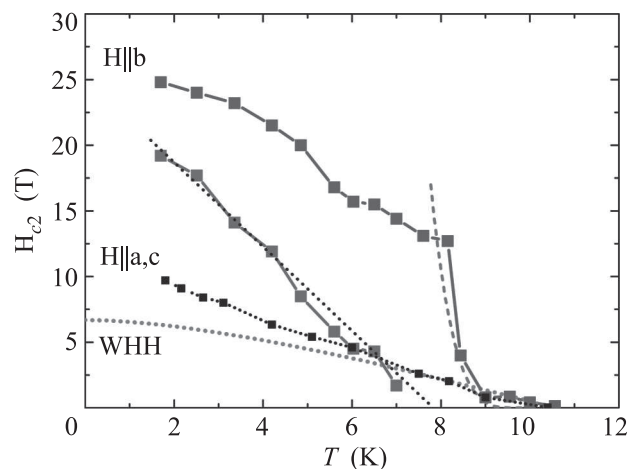


Fig. 1. (Color online) Phase diagram for the k-Br compound with conducting plane oriented parallel (H||ac) and perpendicular (H||b) to the pulsed magnetic fields. Black and blue squares indicate the  $H_{c2}$  values inferred from experiments with H||ac and H||b, respectively. Green squares correspond to the melting transition fields  $H_m(T)$  data. The red dotted line is the best fit of the WHH model to the data for H||ac. Green dotted line is power law fit close to  $T_c$ . [1, 5, 6]. The red dotted line indicate the temperature dependence according to the WHH model neglecting Pauli-limiting

where  $\alpha(T) = \sqrt{2}H_{c2}/H_p(0)$  is the Maki parameter [3]. Whether a superconductor is orbitally limited or Pauli limited can be inferred from the  $\alpha$  parameter value. Above  $H_p$  the Cooper pairs are broken and superconductivity is destroyed. The superconductor made up of weakly-coupled superconducting planes [4] may transform from a three-dimensional system to two-dimensional one [5]. Recently, the change in gra-

dient of  $H_{c2}$  vs  $T$  in [5] was attributed to dimensional crossover from quasi-two-dimensional (low temperatures) to three dimensional (high temperatures). In this short communication, we report (i) results of the magnetic penetration depth measurements in zero-field for  $\kappa$ -(BEDTTTF)<sub>2</sub>Cu[N(CN)<sub>2</sub>]Br single crystals in the temperature range between 1.2 and 15 K and (ii) critical magnetic fields measurements in pulsed fields of up to 56 T in the range 1.2 to 12 K. In both cases, field directions parallel and perpendicular to the conducting ac plane were explored.

Here, we found that below 9 K a definite change in the slope of  $H_{c2}$  vs  $T$  occurs and  $H_{c2}$  begins to follow the power law,  $H_{c2}(T) \propto (9.8 - T^*)^4$ , with fitted  $T^* = 9.8$  K. Notice that this upturn temperature position is almost the same as for  $H_m(T)$ . Similar upturn was observed for  $\lambda$ -(BETS)<sub>2</sub>GaCl<sub>4</sub> organic superconductor, with  $T^*$  as a free parameter. Following [5], we attribute the sharp upturn of  $H_{c2}(T)$  to flux-line lattice melting, and thus characteristic of a two-dimensional superconductor with weakly-coupled layers, because  $H^*$  follows the power like dependence  $H_{c2} \propto (9.8 - T)^4$  dependence, indeed expected from the Gorter-Casimir two-fluid model [6].

Measurements of the temperature dependence of both upper critical magnetic field and melting transition reveal unusual phase diagram for  $\kappa$ -(BEDT-TTF)<sub>2</sub>Cu[N(CN)<sub>2</sub>]Br organic superconductor. The temperature dependence of the upper critical magnetic field  $H_{c2}(T)$  shows an upward divergence as the temperature decreases with a sharp upturn at 9 K, and similar upturn for the temperature dependence of melting field  $H_m(T)$  (see Fig. 1). We tentatively attribute these two upturns to flux-line lattice melting.

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