Flat band and Planckian metal

 $G. E. Volovik^{1)}$

Low Temperature Laboratory, Aalto University, P.O. Box 15100, FI-00076 Aalto, Finland

Landau Institute for Theoretical Physics, 142432 Chernogolovka, Russia

Submitted 20 July 2019 Resubmitted 20 July 2019 Accepted 20 July 2019

DOI: 10.1134/S0370274X19170090

We discuss the recent extension of the Sachdev-Ye-Kitaev (SYK) microscopic model of interacting fermions [1]. The model describes a Planckian metal at low temperatures, in order to explain the linear temperature dependence of their resistivity. We show here that the proposed scenario actually describes the formation of the Khodel-Shaginyan fermion condensate – the finite region of momenta, where the energy of electrons is exactly zero (the flat band) [2]. The microscopic derivation of the flat band in this interacting model supports the original idea of Khodel and Shaginyan based on the phenomenological approach. It also suggests that it is the flat band, which is responsible for the linear dependence of resistivity on temperature in "strange metals".

There are different potential sources of the formation of the electronic flat band with zero energy, see, e.g., [3]. In particular it can be formed due to electronelectron interaction. The flat band formed by interaction has been first discussed by Khodel and Shaginyan (KS) in 1990 [2], who used the phenomenological Landau theory of Fermi liquid, see also [4–7] and Fig. 1. This dispersionless energy spectrum has a singular density of states. As a result the superconducting gap and transition temperature are proportional to the coupling constant instead of the exponential suppression in conventional metals with Fermi surfaces. For nuclear systems the linear dependence of the superconducting gap on the coupling constant has been found by Belyaev [8]. In a more rigourous manner the flat band induced by interaction has been obtained in [9, 10]. Experimentally the merging of levels at the Fermi surface due to interaction has been reported in [11, 12]

In twisted bilayer graphene there is indication that interaction leads to the further flattening of the spectrum [13, 14] in addition to the geometrical/topological flattening caused by the magic angle twist [15, 16].

¹⁾e-mail: volovik@boojum.hut.fi

 $\delta E\{n(p)\} = \int \varepsilon(p) \,\delta n(p) d^d p = 0$ Two solutions: $\delta n(p) = 0$ & $\varepsilon(p) = 0$

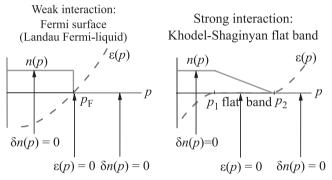


Fig. 1. (Color online) Consequences of Landau theory of Fermi liquid. Variation of the energy functional over the occupancy $n(\mathbf{p})$ gives two possible solutions: $\epsilon(\mathbf{p}) = 0$ and $\delta n(\mathbf{p}) = 0$ $(n(\mathbf{p}) = 0$ or $n(\mathbf{p}) = 1$). Left: The Landau Fermi liquid, where the solution $\epsilon(\mathbf{p}) = 0$ takes place on Fermi surface. Outside of Fermi surface one has either $n(\mathbf{p}) = 0$ or $n(\mathbf{p}) = 1$. Right: Khodel–Shaginyan flat band, where the solution $\epsilon(\mathbf{p}) = 0$ takes place in the finite region of momentum space. In this region $0 < n(\mathbf{p}) < 1$

In recent paper by Patel and Sachdev [1] the lattice extension of the Sachdev–Ye–Kitaev (SYK) model has been used to study the problem of the "bad metal" with the universal linear dependence of resistivity on temperature [17–20]. However, it appears that signatures of the KS flat band in Fig. 1 (right panel) are very similar to those in Figs. 2a and 3a from Patel and Sachdev (PS) [1]. Indeed, Fig. 2a from the PS paper shows the occupancy $n(\mathbf{p})$, which exhibits the same behavior as $n(\mathbf{p})$ in Fig. 1 (right panel), with the finite region where $0 < n(\mathbf{p}) < 1$. According to Khodel–Shaginyan, in this region the quasiparticle energy should be zero. And this is clearly seen from the electron spectral density shown in Fig. 3a of the PS paper. So one may conclude that the extended SYK model provides another possible realization of the KS flat band.

That is why the extended SYK model can be used for studying different properties of the materials which experience formation of the KS flat band, including possibly the "bad metal" behavior. In this model, the universal linear dependence of resistivity on temperature has been obtained [1] in the regime, where the signatures of the flat band are transparent. From that one may conclude that the phenomenon of Planck metal or bad metal is the consequence of the Khodel–Shaginyan flat band emerging in this model. The idea that the flat band may serve as the origin of the "strange metal" behavior has been suggested earlier, see, e.g., [21], and recent papers [22, 23].

This work has been supported by the European Research Council (ERC) under the European Union's Horizon 2020 research and innovation programme (Grant Agreement # 694248).

Full text of the paper is published in JETP Letters journal. DOI: 10.1134/S002136401917003X

- A. A. Patel and S. Sachdev, Phys. Rev. Lett. 123, 066601 (2019).
- V. A. Khodel and V. R. Shaginyan, JETP Lett. 51, 553 (1990).
- G. E. Volovik, Pis'ma ZhETF 107, 537 (2018) [JETP Lett. 107, 516 (2018)].
- 4. G.E. Volovik, JETP Lett. 53, 222 (1991).
- 5. P. Nozieres, J. Phys. (Fr.) **2**, 443 (1992).
- 6. G.E. Volovik, JETP Lett. 59, 830 (1994).
- T. T. Heikkilä and G. E. Volovik, in *Basic Physics of Functionalized Graphite*, Springer (2016), p. 123.
- 8. S.T. Belyaev, JETP 12, 968 (1961).
- D. Yudin, D. Hirschmeier, H. Hafermann, O. Eriksson, A. I. Lichtenstein, and M. I. Katsnelson, Phys. Rev. Lett. **112**, 070403 (2014).

- 10. S.-S. Lee Phys. Rev. D **79**, 086006 (2009).
- A. A. Shashkin, V. T. Dolgopolov, J. W. Clark, V. R. Shaginyan, M. V. Zverev, and V. A. Khodel, Phys. Rev. Lett. **112**, 186402 (2014).
- M. Yu. Melnikov, A.A. Shashkin, V.T. Dolgopolov, S.-H. Huang, C.W. Liu, and S. V. Kravchenko, Sci. Rep. 7, 14539 (2017).
- D. Marchenko, D.V. Evtushinsky, E. Golias, A. Varykhalov, Th. Seyller, and O. Rader, Sci. Adv. 4, 0059 (2018).
- S. Carr, Sh. Fang, P. Jarillo-Herrero, and E. Kaxiras, Phys. Rev. B 98, 085144 (2018).
- Y. Cao, V. Fatemi, Sh. Fang, K. Watanabe, T. Taniguchi, E. Kaxiras, and P. Jarillo-Herrero, Nature 556, 43 (2018).
- Y. Cao, V. Fatemi, A. Demir, Sh. Fang, S. L. Tomarken, J. Y. Luo, J. D. Sanchez-Yamagishi, K. Watanabe, T. Taniguchi, E. Kaxiras, R. C. Ashoori, and P. Jarillo-Herrero, Nature 556, 80 (2018).
- A. Legros, S. Benhabib, W. Tabis et al. (Collaboration), Nat. Phys. 15, 142 (2018).
- Y. Nakajima, T. Metz, C. Eckberg et al. (Collaboration), arXiv:1902.01034.
- Y. Cao, D. Chowdhury, D. Rodan-Legrain, O. Rubies-Bigorda, K. Watanabe, T. Taniguchi, T. Senthil, and P. Jarillo-Herrero, arXiv:1901.03710.
- P. T. Brown, D. Mitra, E. Guardado-Sanchez, R. Nourafkan, A. Reymbaut, C.-D. Hebert, S. Bergeron, A. M. S. Tremblay, J. Kokalj, D. A. Huse, P. Schau, and W. S. Bakr, Science **363**, 379 (2019).
- V. R. Shaginyan, K. G. Popov, and V. A. Khodel, Phys. Rev. B 88, 115103 (2013).
- V. A. Khodel, J. W. Clark, and M. V. Zverev, Phys. Lett. A 382, 3281 (2018).
- V. R. Shaginyan, M. Ya. Amusia, A. Z. Msezane, V. A. Stephanovich, G. S. Japaridze, and S. A. Artamonov, Pis'ma v ZhETF 109, 266 (2016)